

Couplage magnéto-électrique dans les matériaux multiferroïques

Renaud Schleck

LPEM-ESPCI

Ricardo Lobo

LPEM-ESPCI-CNRS

M.B. Lepelet, J.Varignon

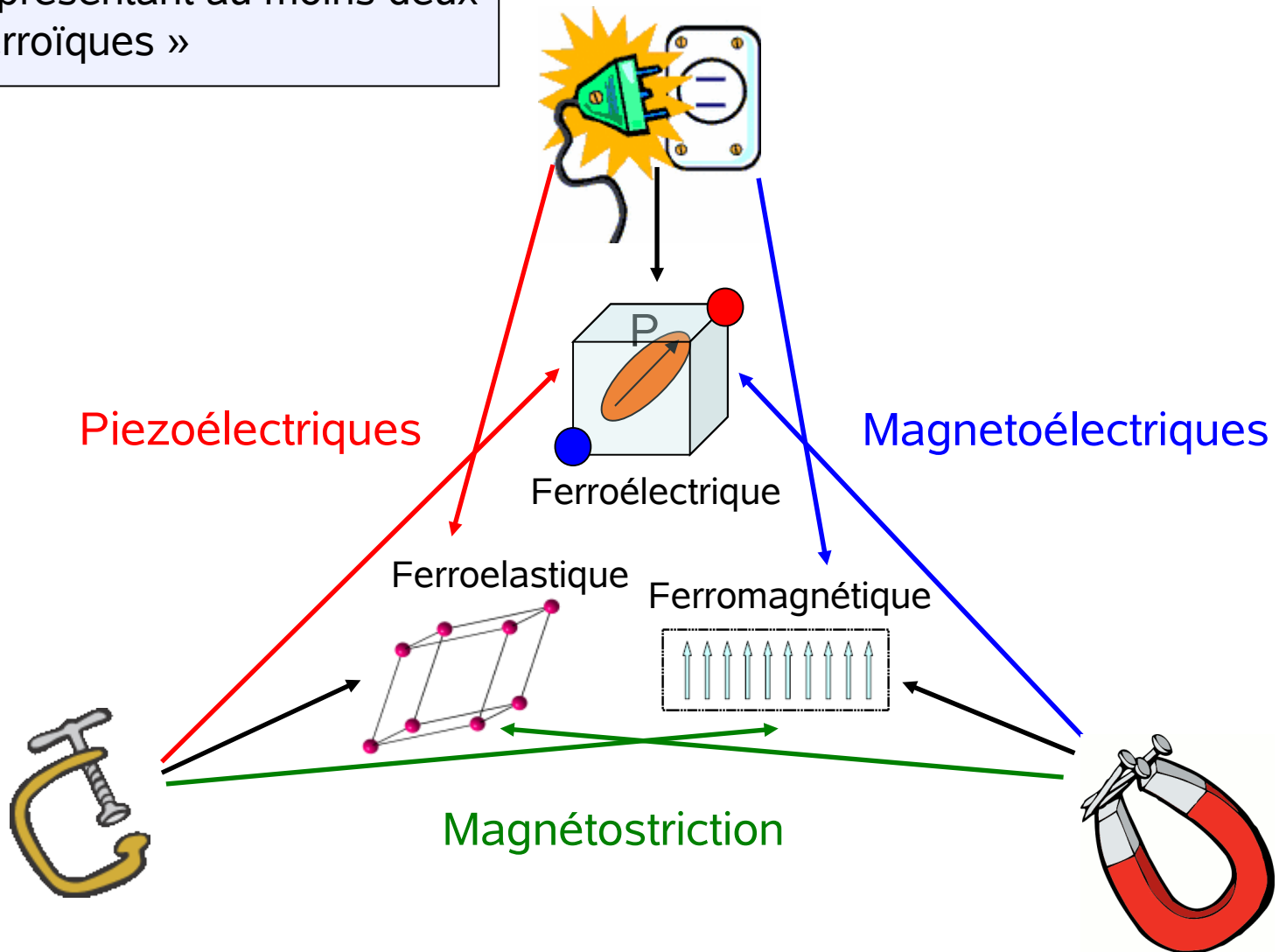
CRISMAT-ENSICAEN

H. Sakata

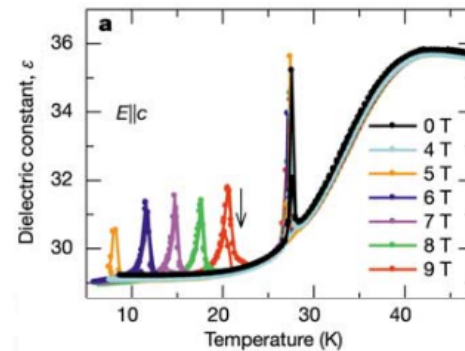
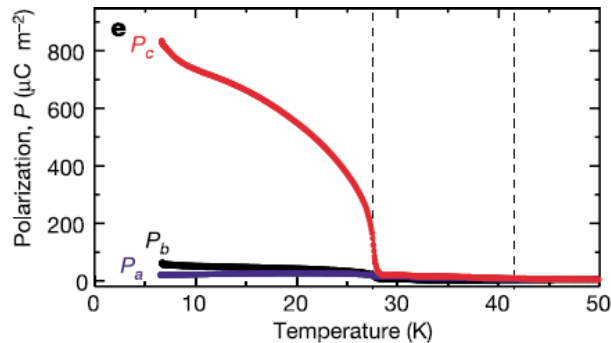
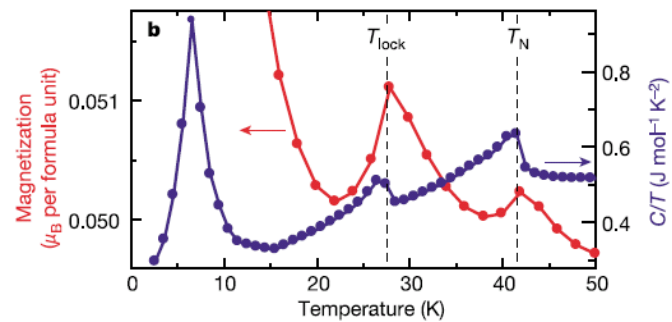
Tokyo University of Technology

Les matériaux multiferroïques

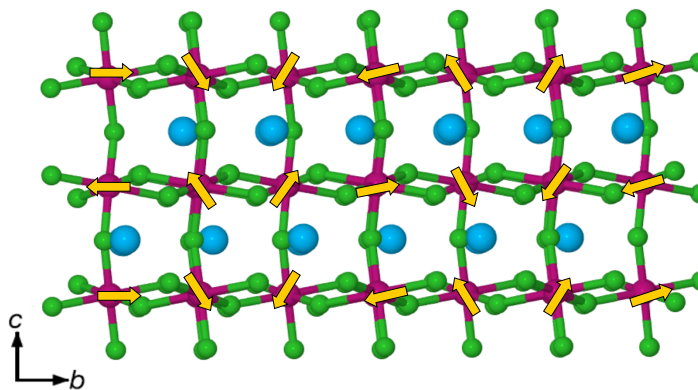
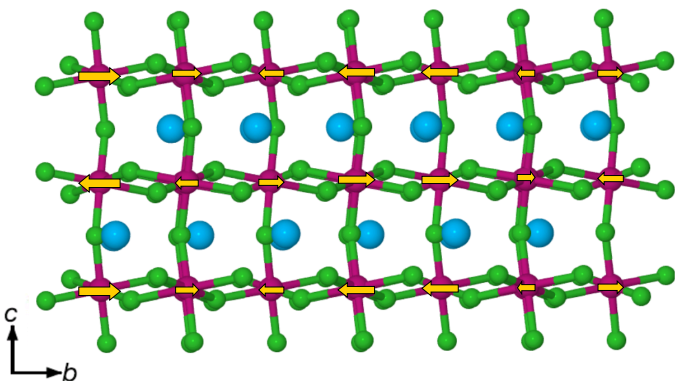
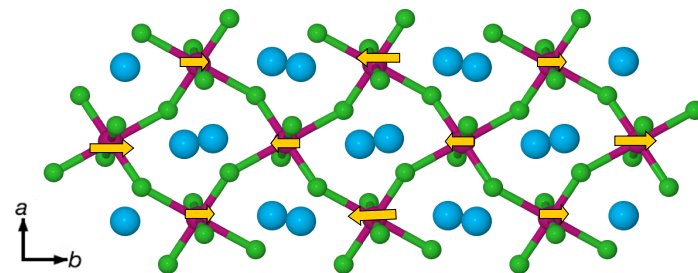
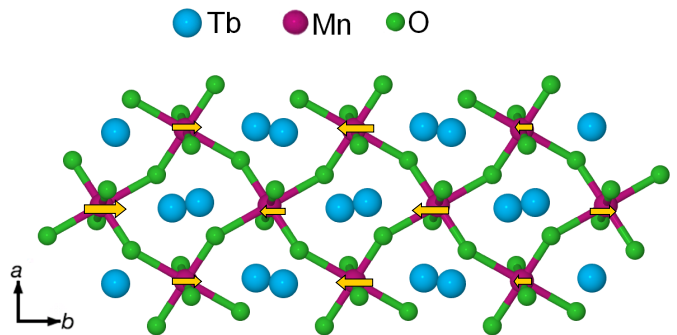
Matériaux présentant au moins deux ordres « ferroïques »



TbMnO₃ : Structure et transitions de phase



Kimura et al., *Nature* 426, 58 (2003)



$$T_C < T < T_N$$

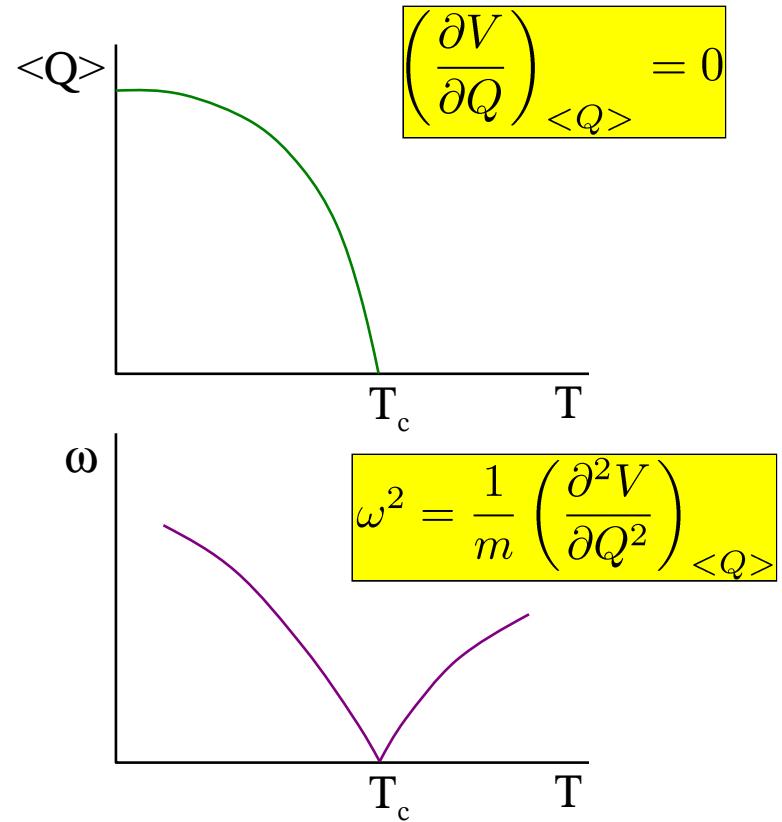
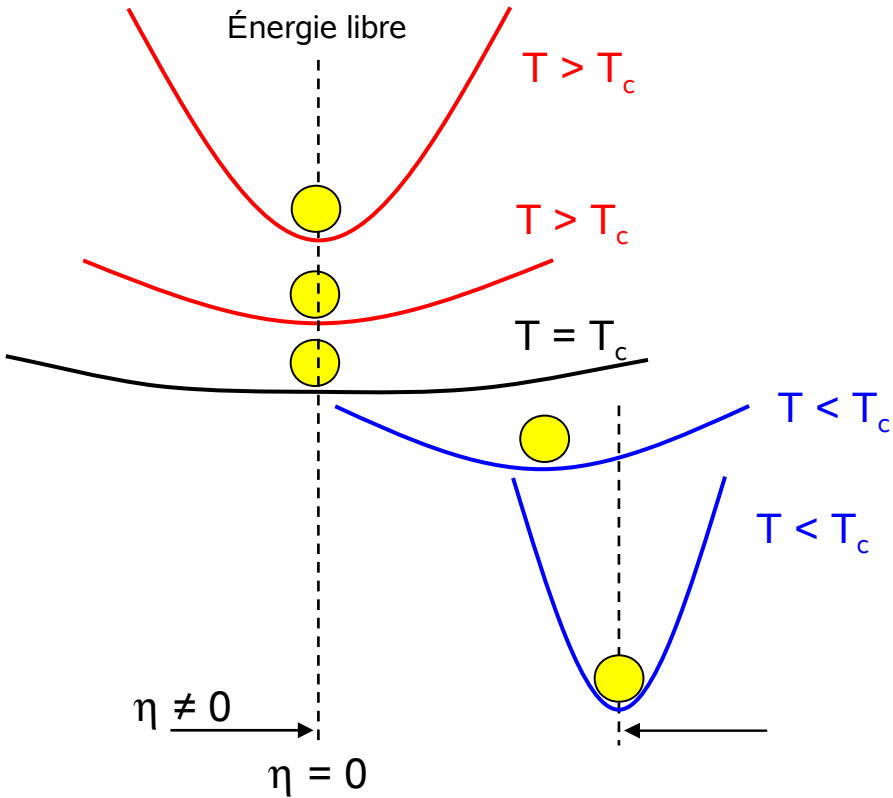
$$T_N = 41 \text{ K} \quad T_C = 27 \text{ K}$$

$$T < T_C$$

Ferroélectricité et mode mou

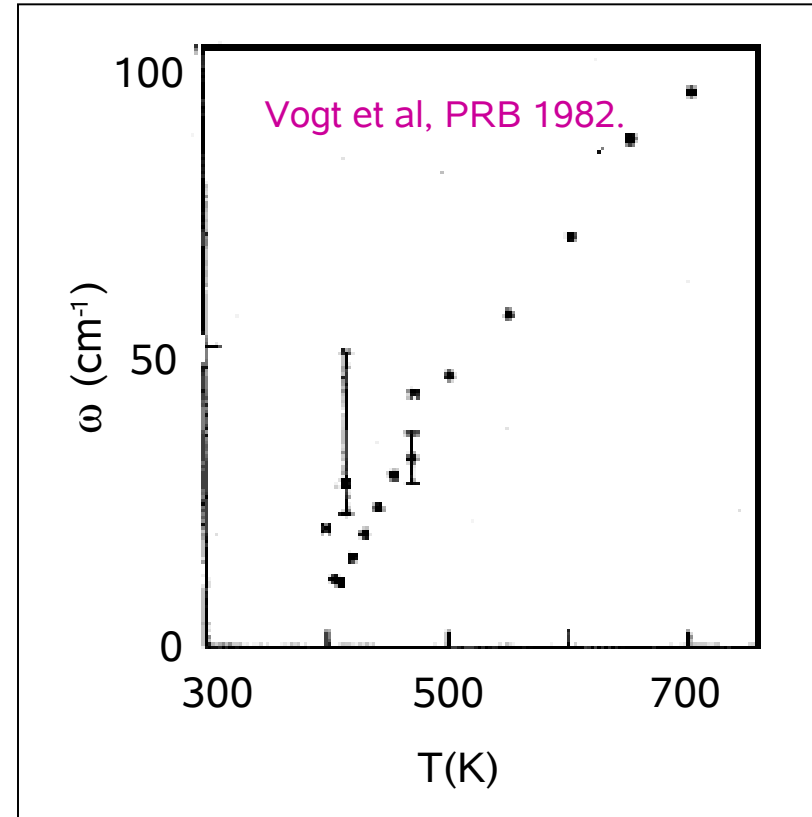
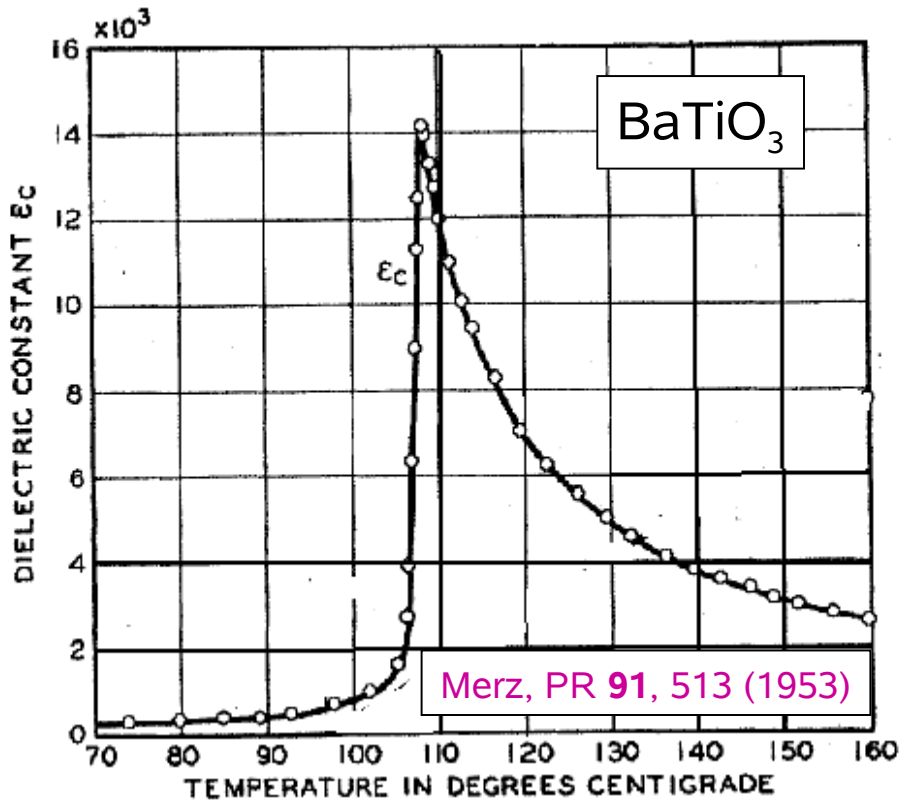
Q est une variable dynamique du système

Parmètre d'ordre $\eta = \langle Q \rangle$ $\begin{cases} = 0, T > T_c \\ \neq 0, T < T_c \end{cases}$

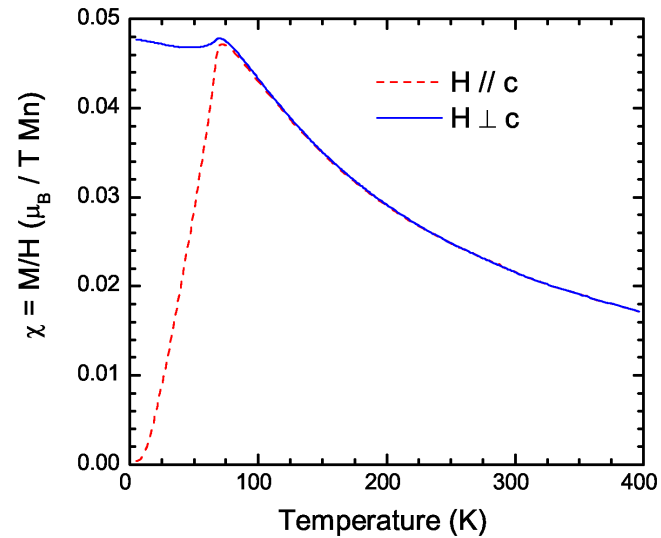
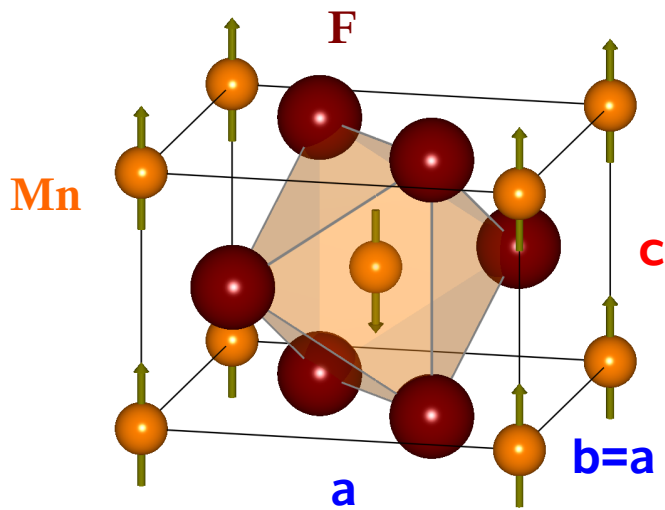


Ferroélectricité et mode mou

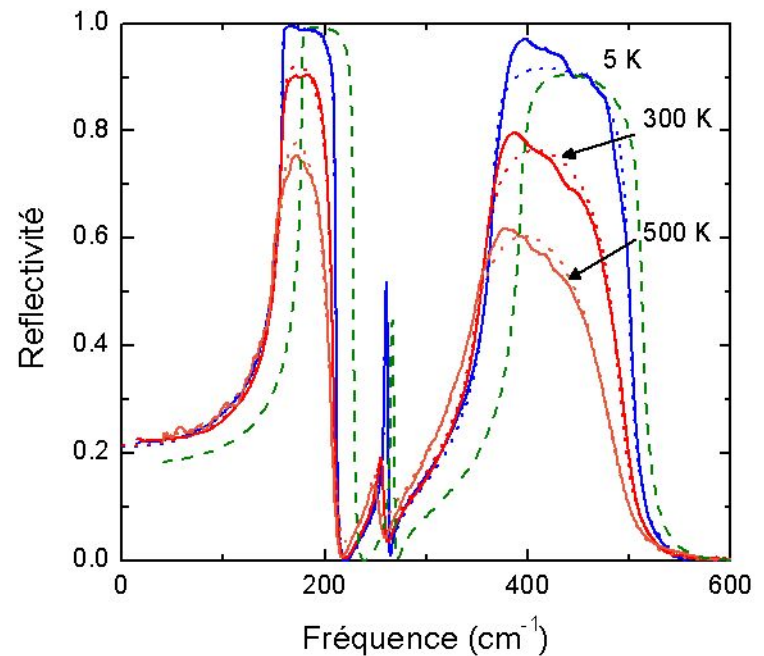
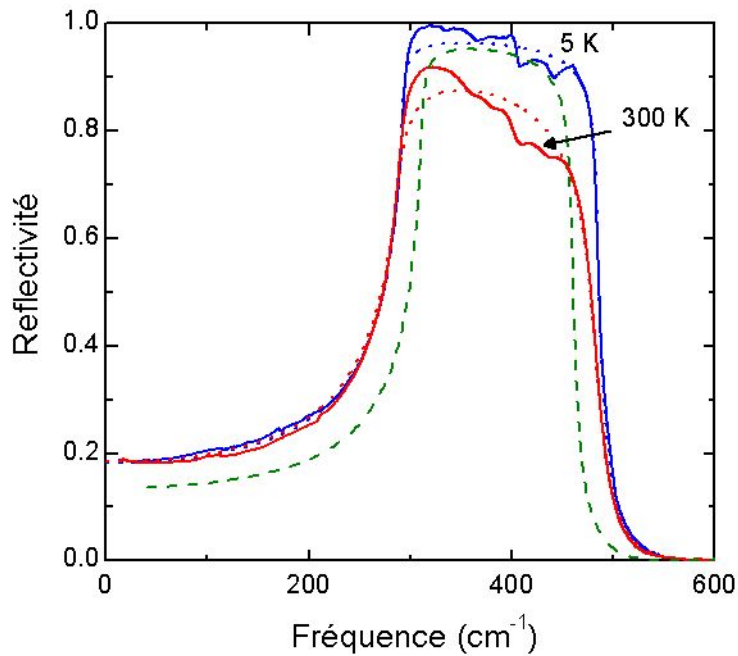
Pour un phonon isolé :
$$\frac{2}{\pi} \int \sigma_1(\omega) d\omega = \Delta \epsilon_j \Omega_{0j}^2 \propto \frac{n_j q_j^2}{m_j} = \text{const}$$



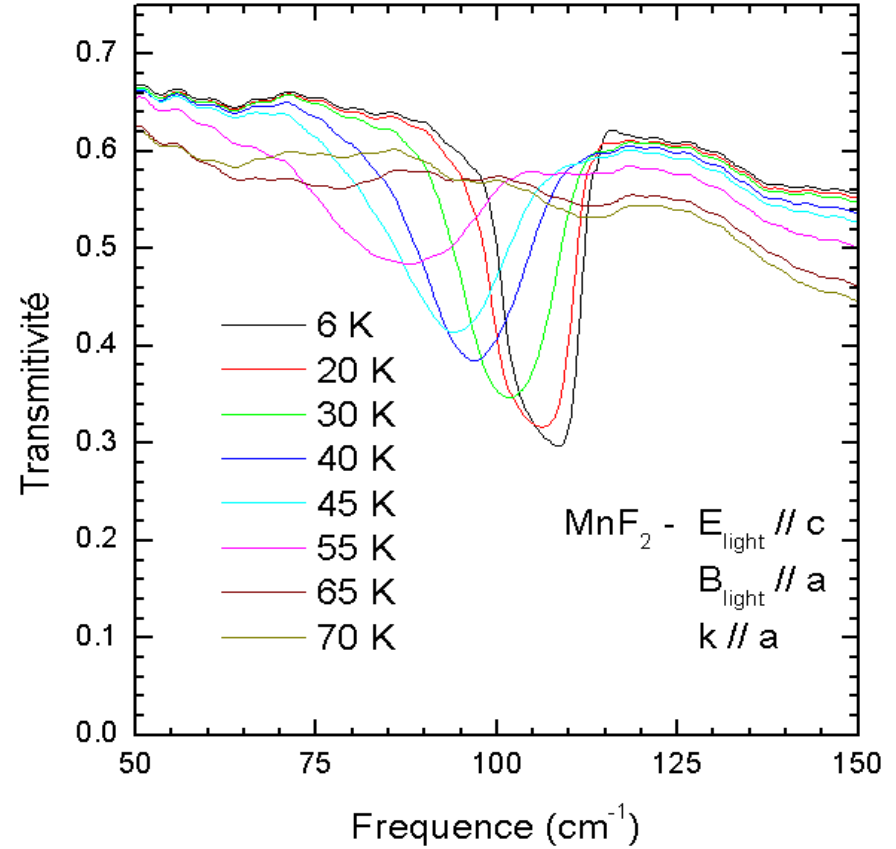
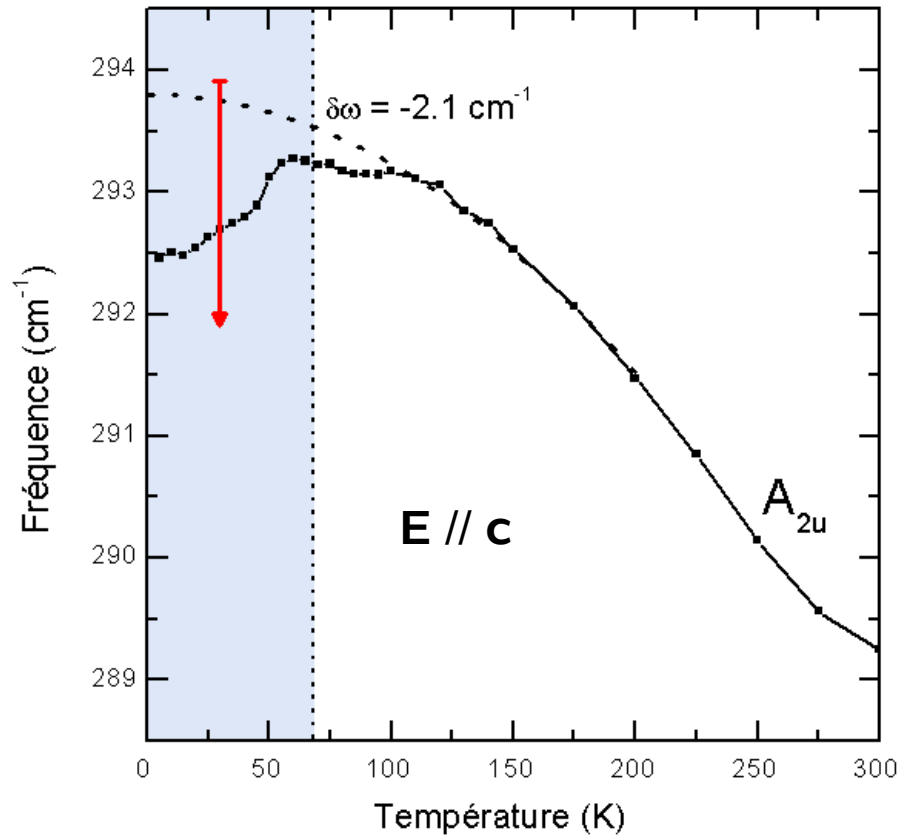
Plus simple : MnF_2 , simplement antiferromagnétique



$T_N = 67 \text{ K}$



MnF₂ : réponse infrarouge de l'axe c



MnF₂ : couplage magnéto-électrique : règle de somme

- Règle de somme optique (conservation du nombre de charges)

$$\int_0^{\infty} \text{Re}[\sigma(\omega)] d\omega = \text{const} \times \frac{n_e}{m_e}$$

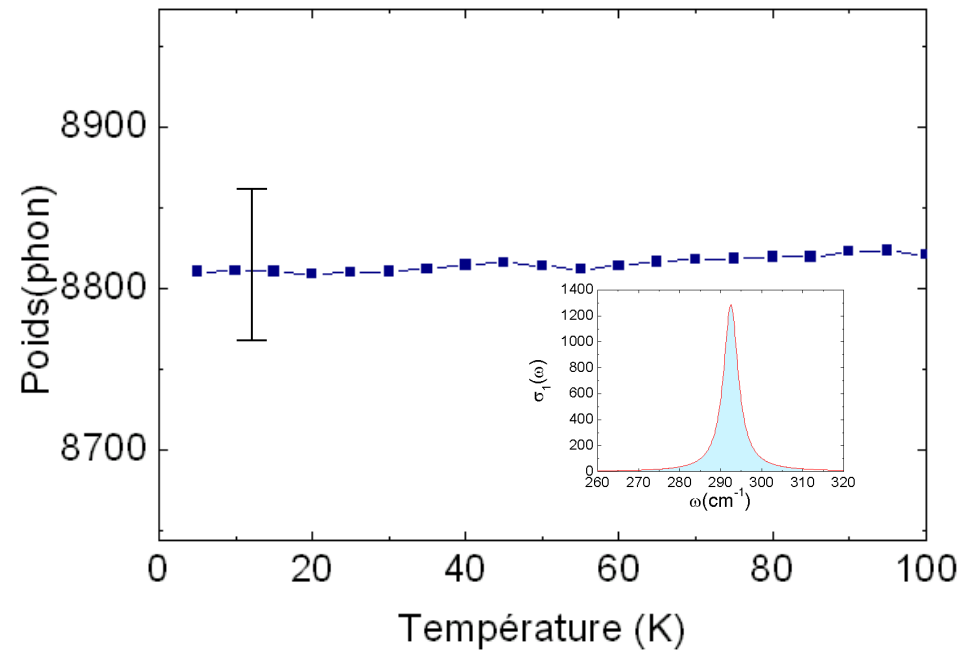
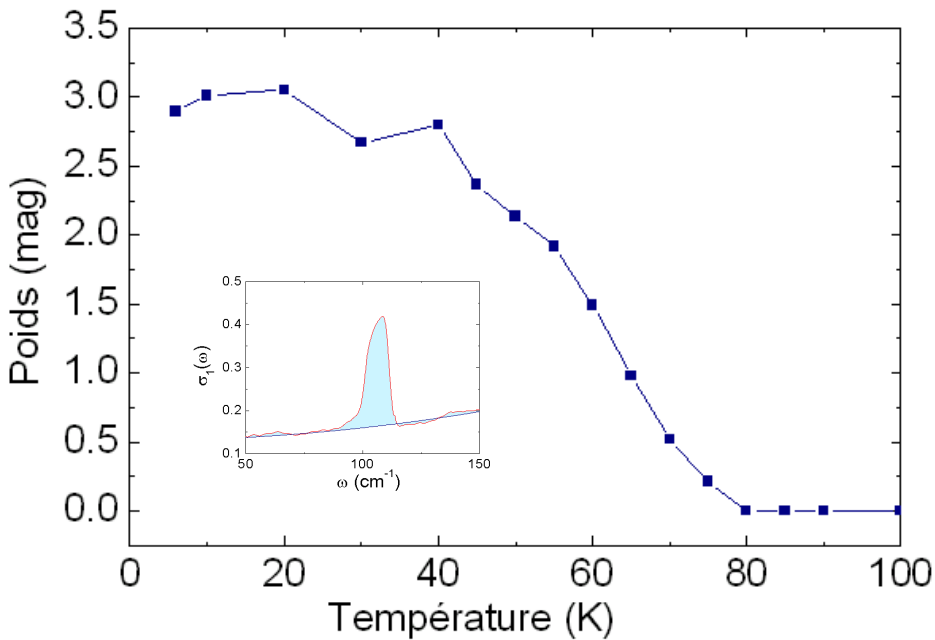
- Pour une excitation couplée à aucune autre :

$$\int_{\text{Excitation}} \sigma_1(\omega) d\omega = \text{const}$$

- Si le système {magnons, phonons} est indépendant :

$$\underbrace{\int_{\text{Magnons}} \sigma_1(\omega) d\omega}_{\text{Poids spectral des magnons}} + \underbrace{\int_{\text{Phonons}} \sigma_1(\omega) d\omega}_{\text{Poids spectral des phonons}} = \text{const}$$

MnF₂ : couplage magnéto-électrique : règle de somme



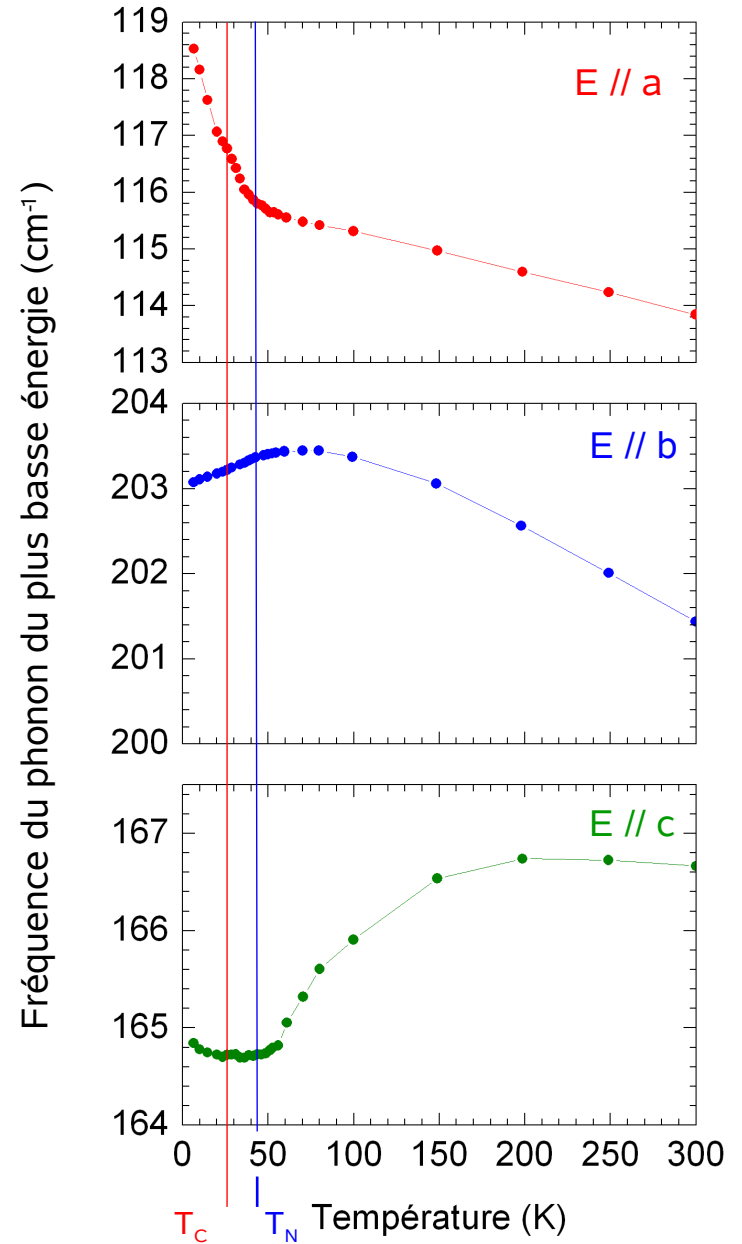
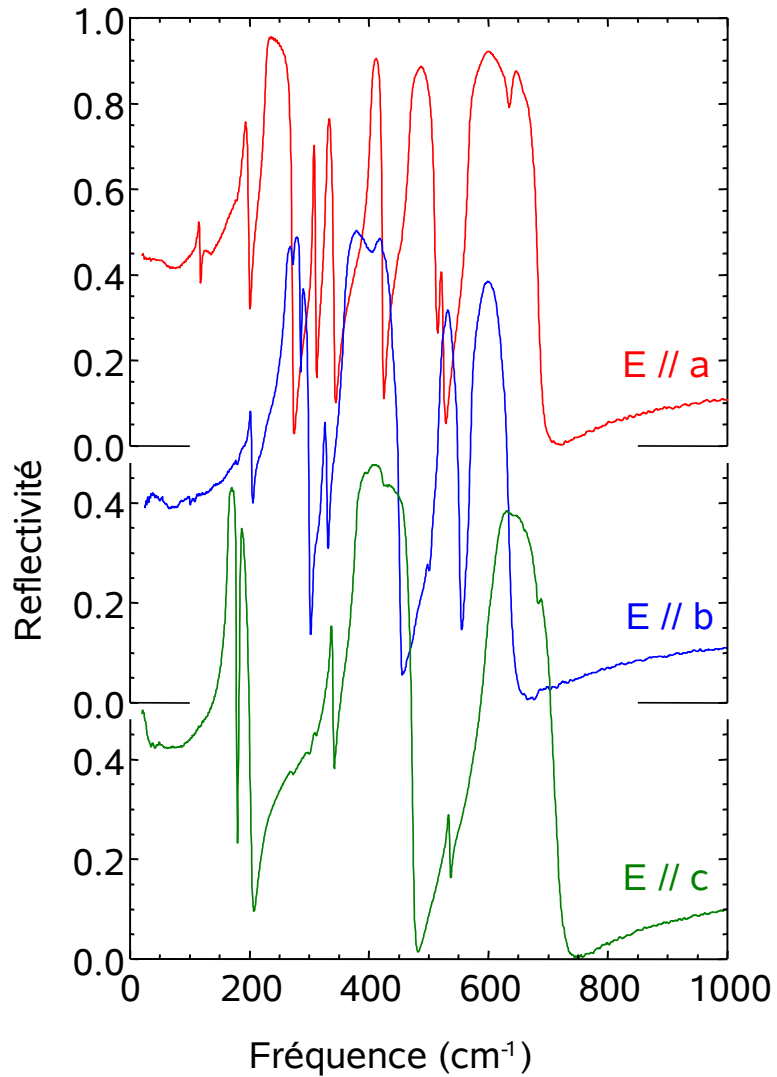
$$Poids(mag) = \int_{2-Magnon} \sigma_1(\omega) d\omega \approx 2.8$$

$$Poids(phon) = \int_{Phonon} \sigma_1(\omega) d\omega \approx 8400$$

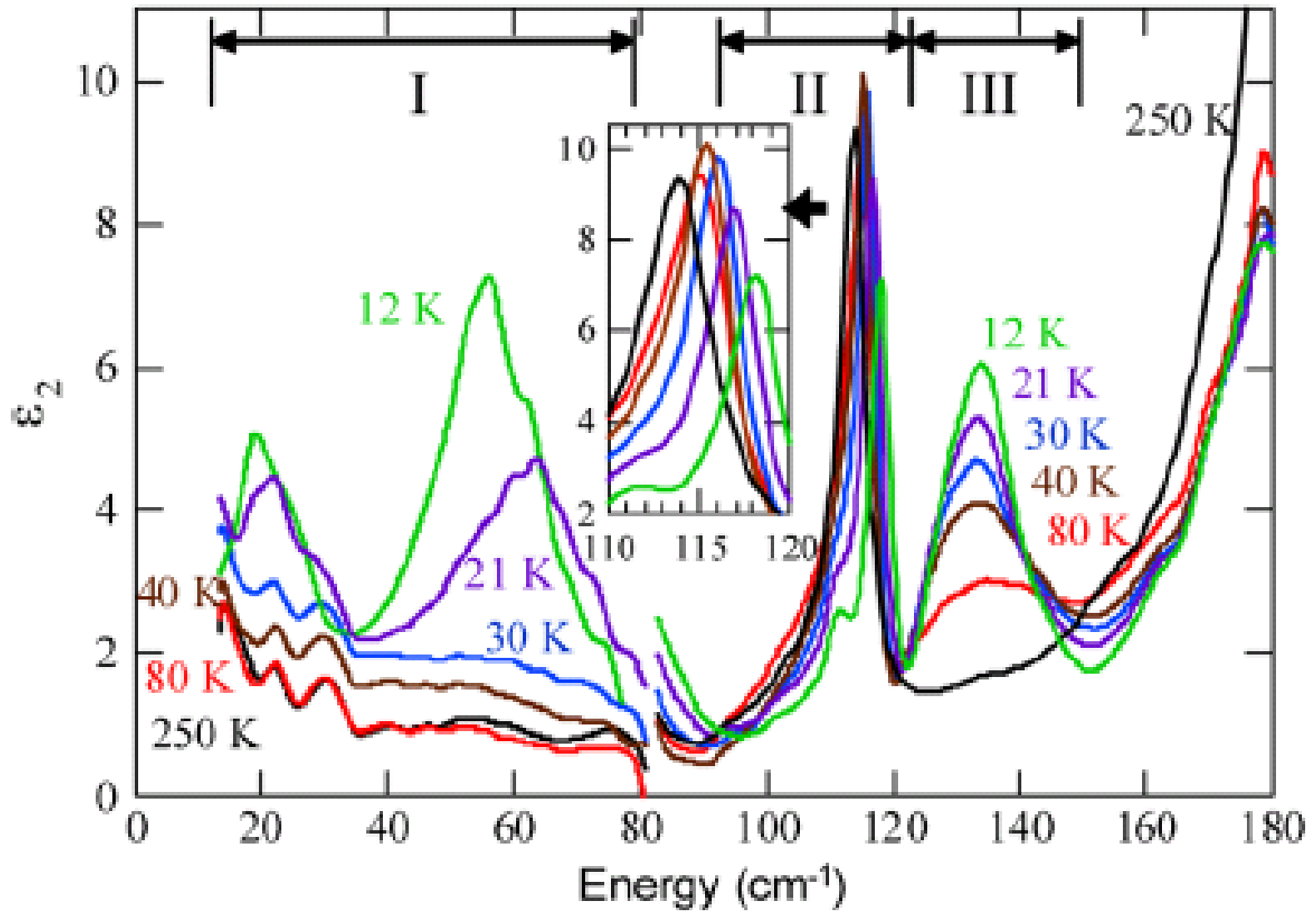
$$\frac{Poids(mag)}{Poids(phon)} \approx 3 \cdot 10^{-4}$$

On ne peut pas voir un changement relatif de $3 \cdot 10^{-4}$ dans le poids spectral du phonon ! Donc on ne peut pas conclure sur un éventuel transfert de poids spectral.

TbMnO₃ : Reflectivité

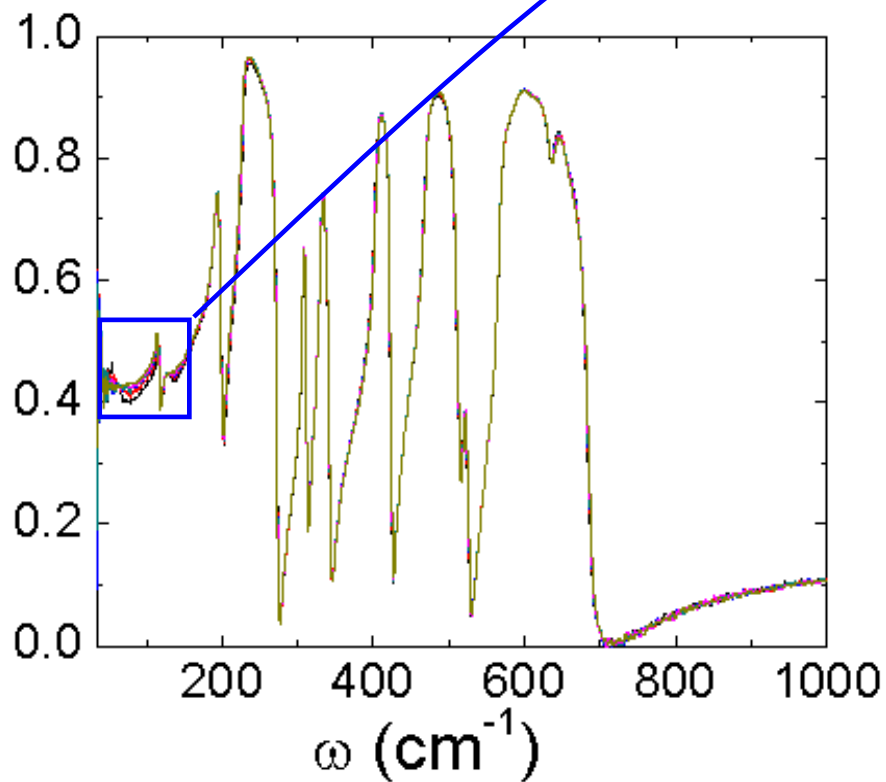


TbMnO₃ : l'axe a

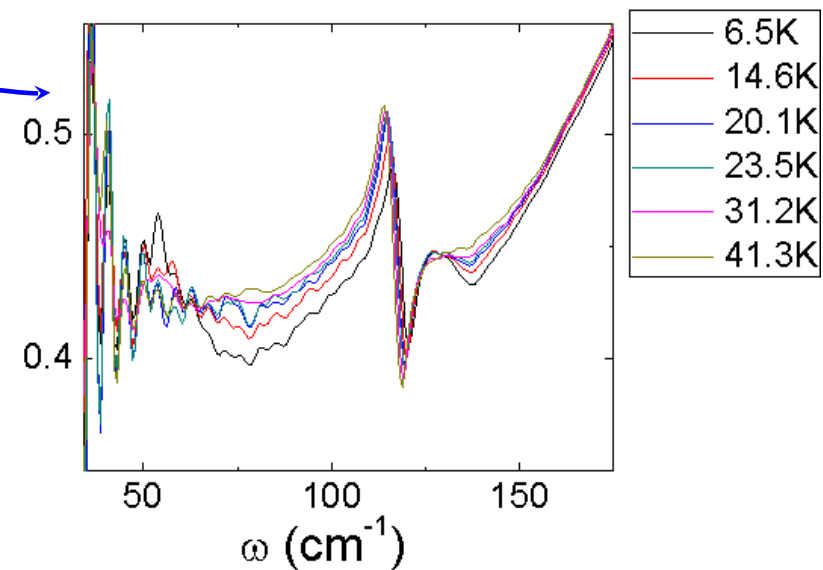


TbMnO₃ : l'axe a

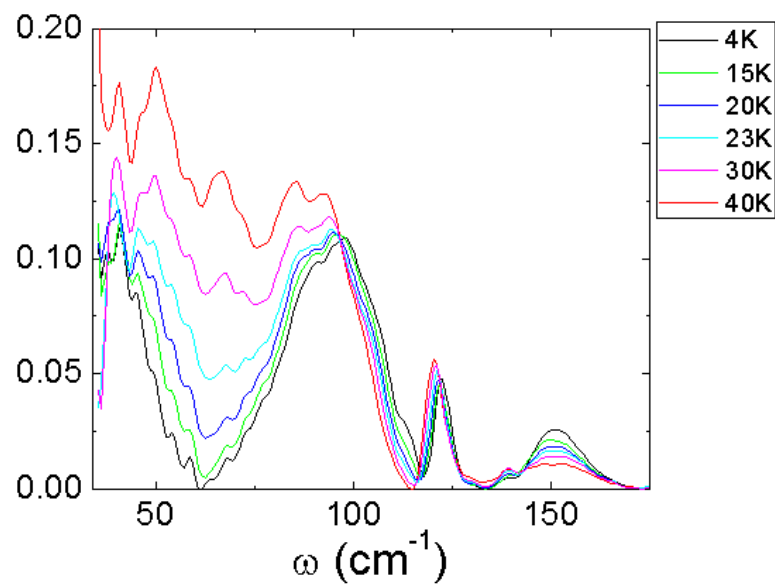
Reflectivité



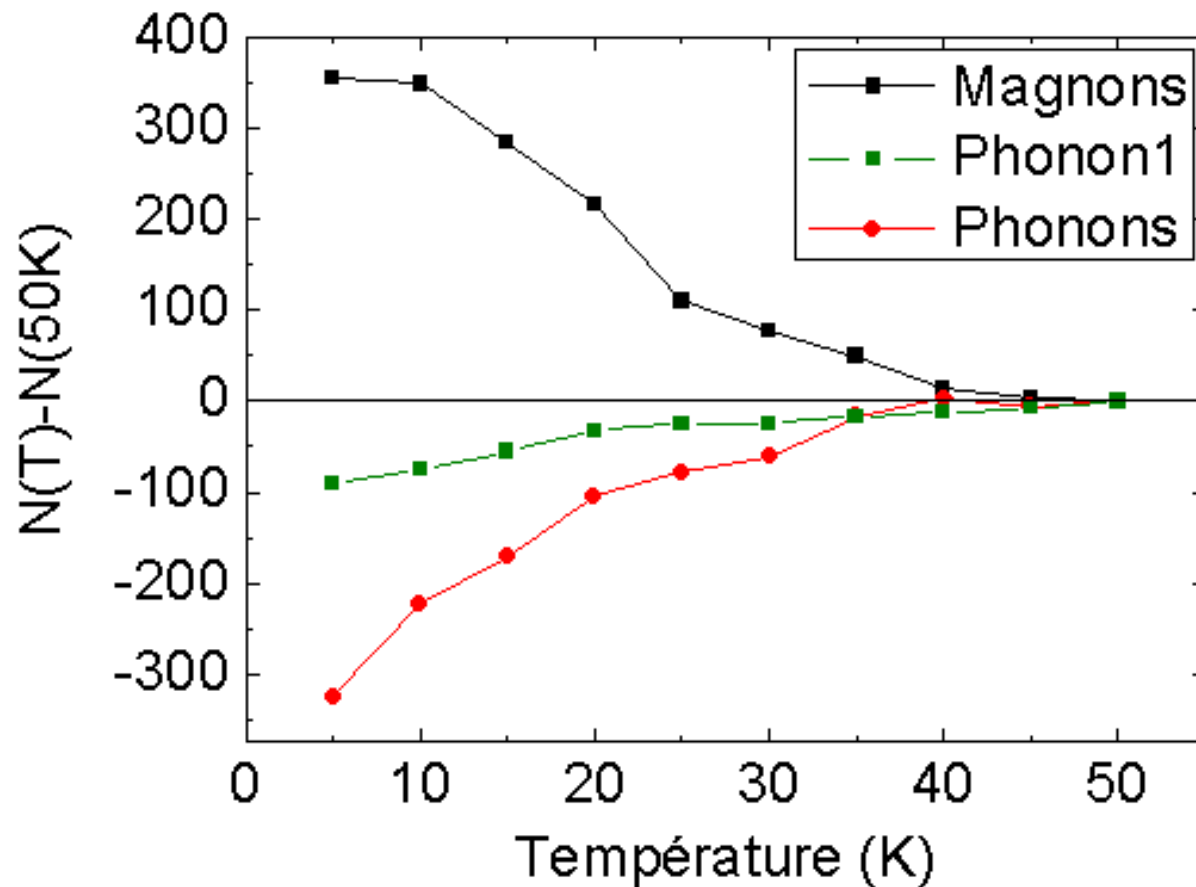
Reflectivité



Transmission

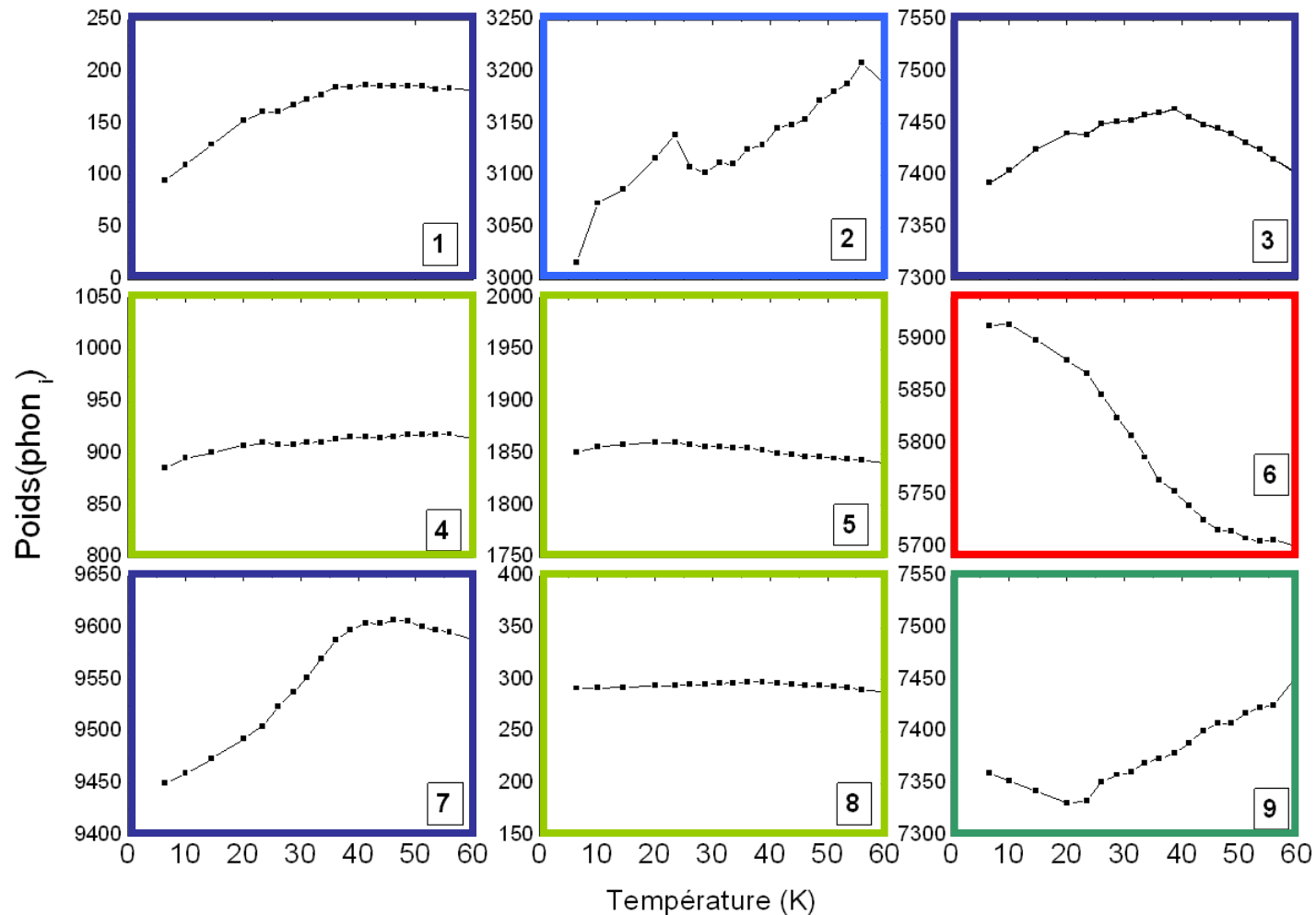


TbMnO₃ : transfert de poids spectral



Il y a un transfert de poids spectral des magnons vers **tous** les phonons quand la température augmente .

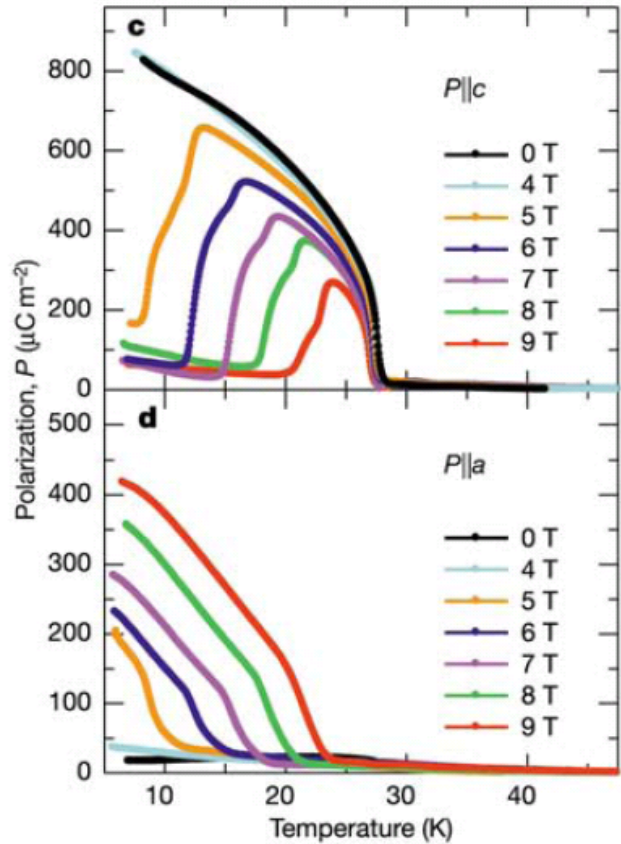
TbMnO₃ : transfert de poids spectral... vers quels phonons ?



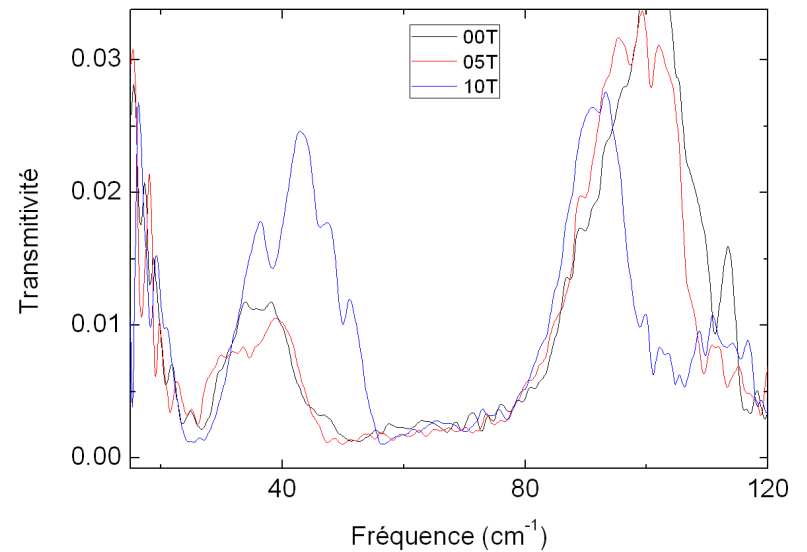
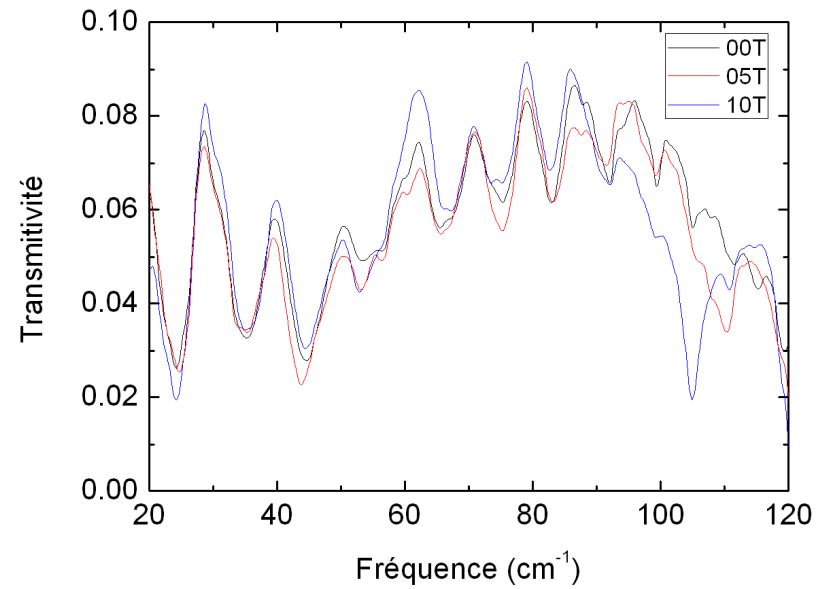
A priori, les phonons 1,2,3,6,7 et 9 sont couplés aux magnons.

Mais 6 et 7 forment probablement un système isolé.

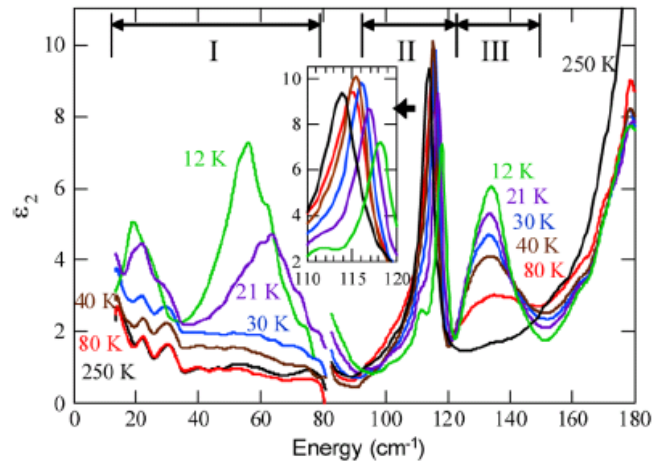
TbMnO₃ : et avec un champ magnétique ?



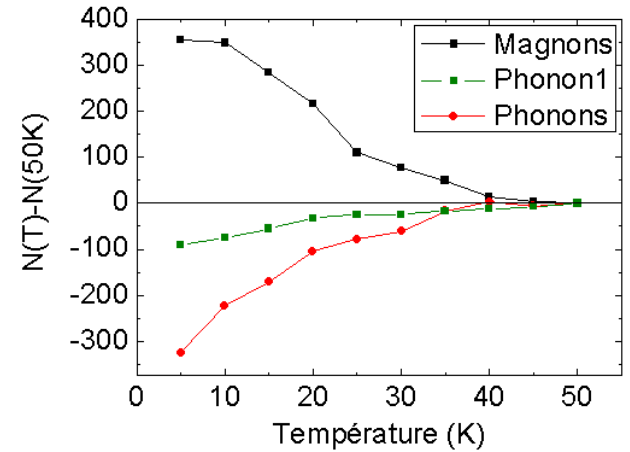
Kimura et al., *Nature* 426, 58 (2003)



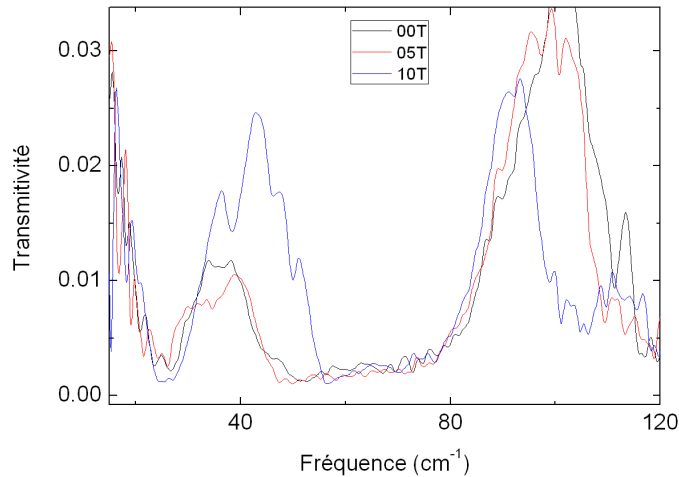
Conclusion



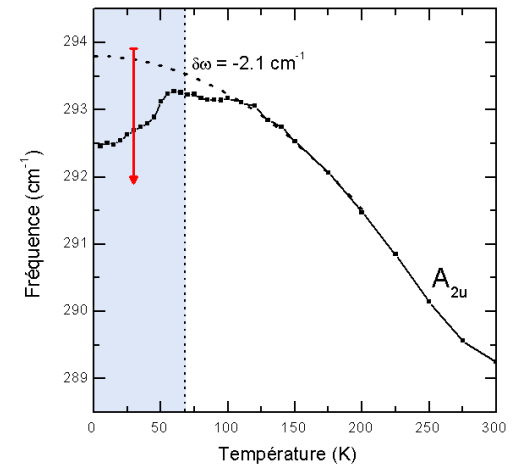
- Dans TbMnO_3 il y a un transfert de poids spectral entre les excitations magnétiques et les phonons



- Le transfert s'effectue vers l'ensemble des phonons.

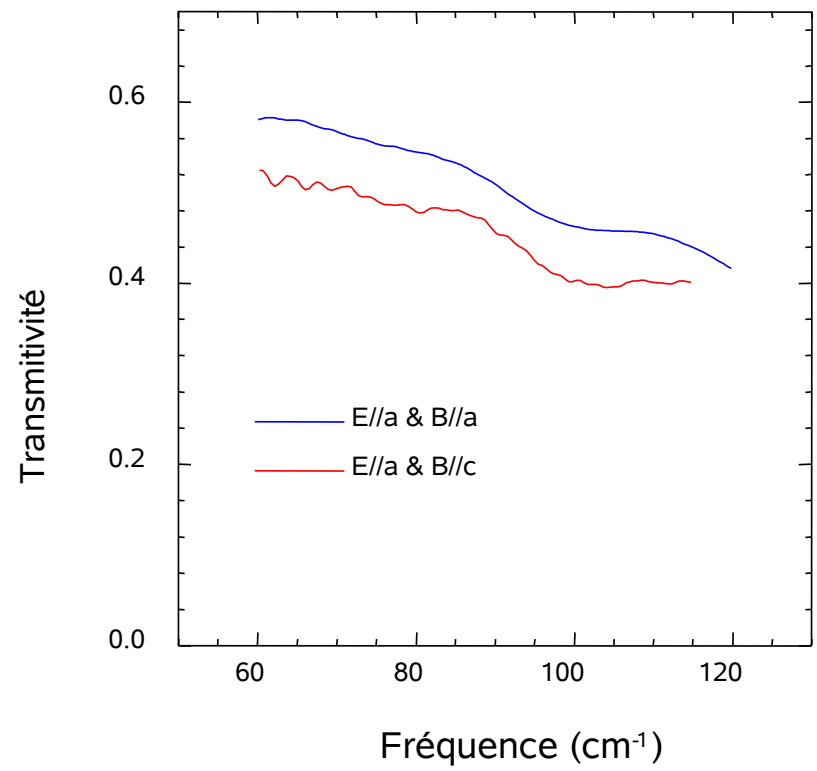
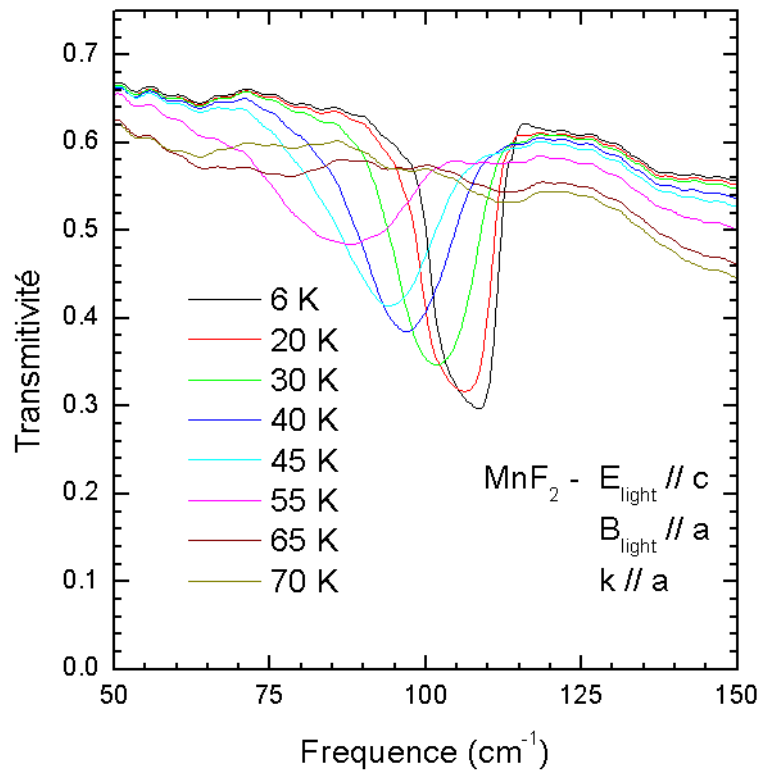


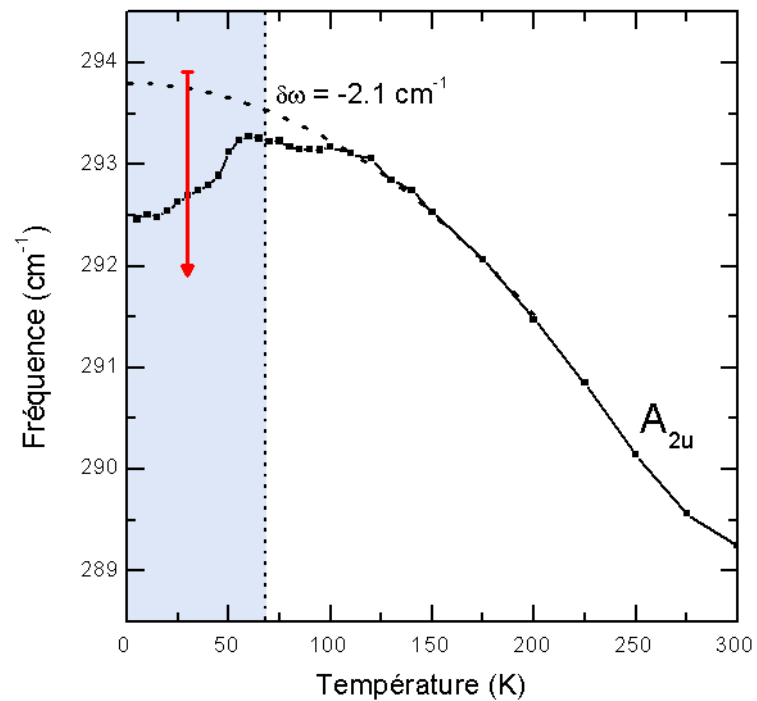
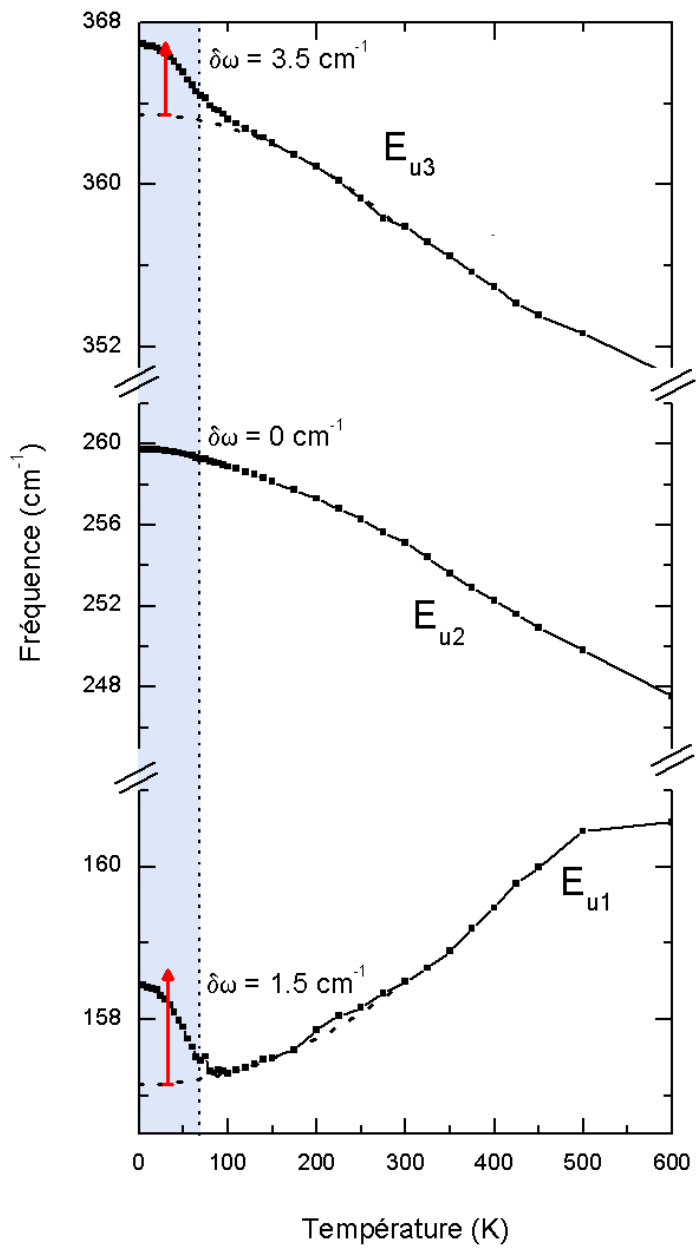
- L'application d'un champ magnétique selon **b** détruit une partie de l'excitation magnétique suivant **a**



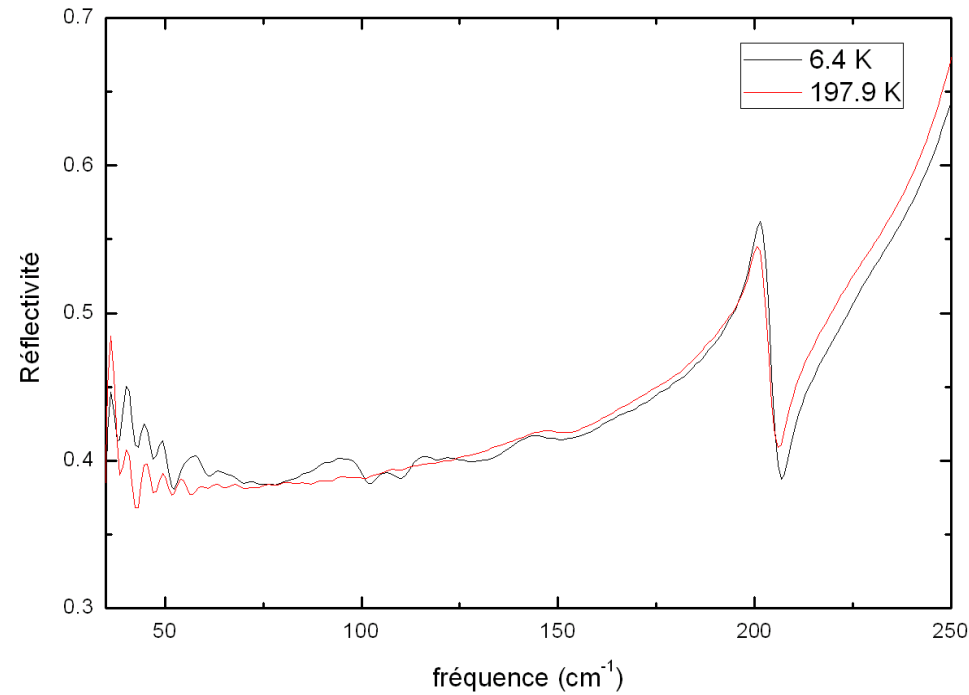
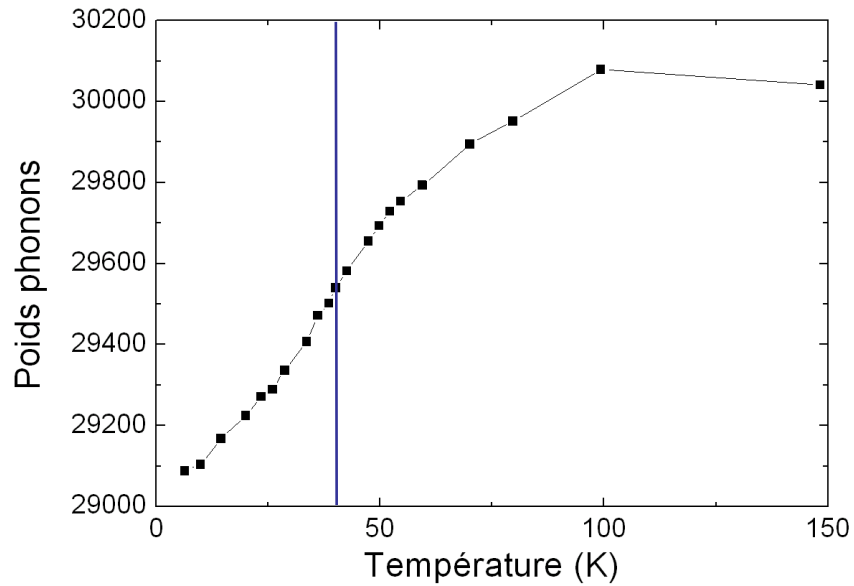
- Même les systèmes antiferromagnétiques les plus simples présentent un couplage magnéto-électrique (cf MnF_2)

FIN



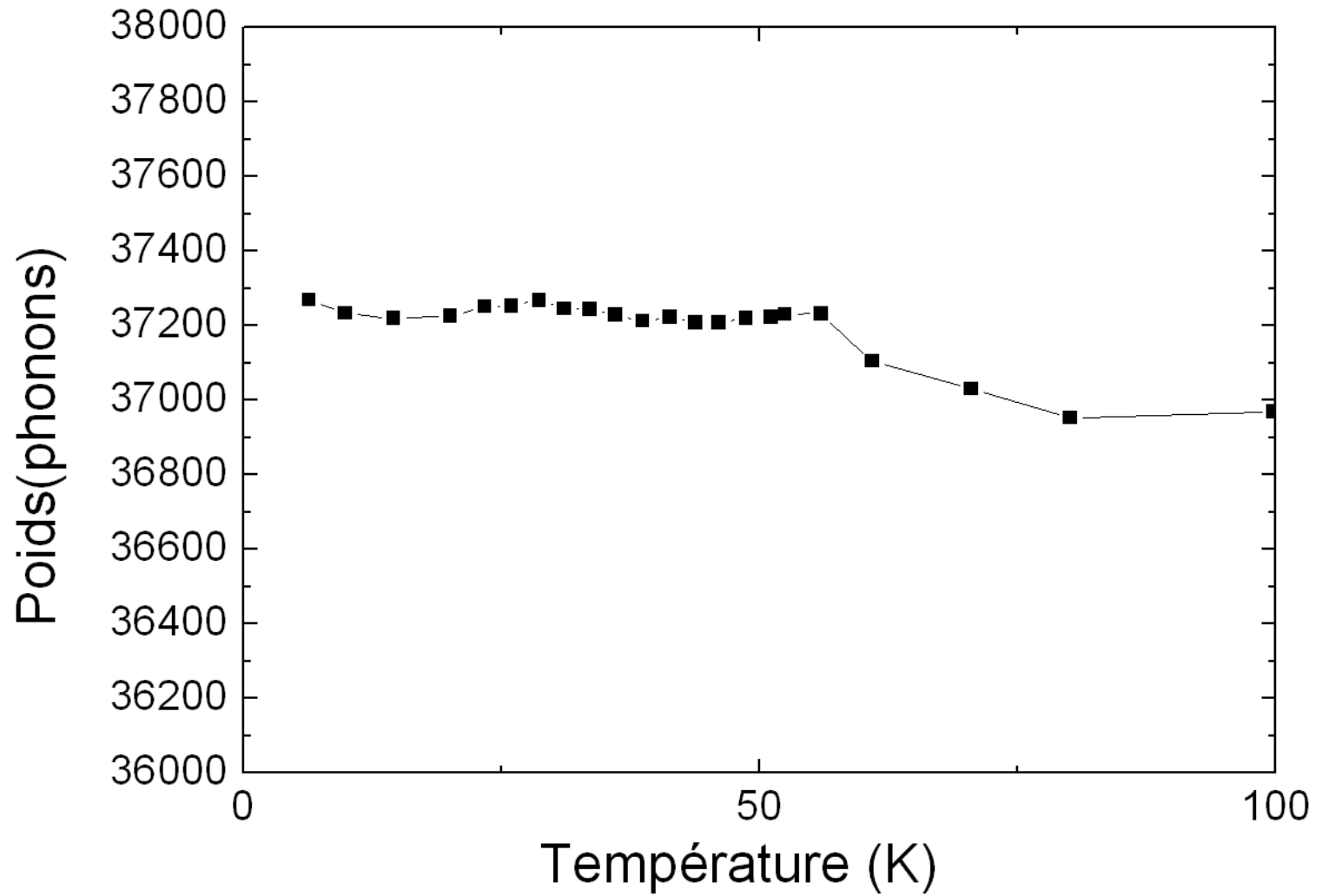


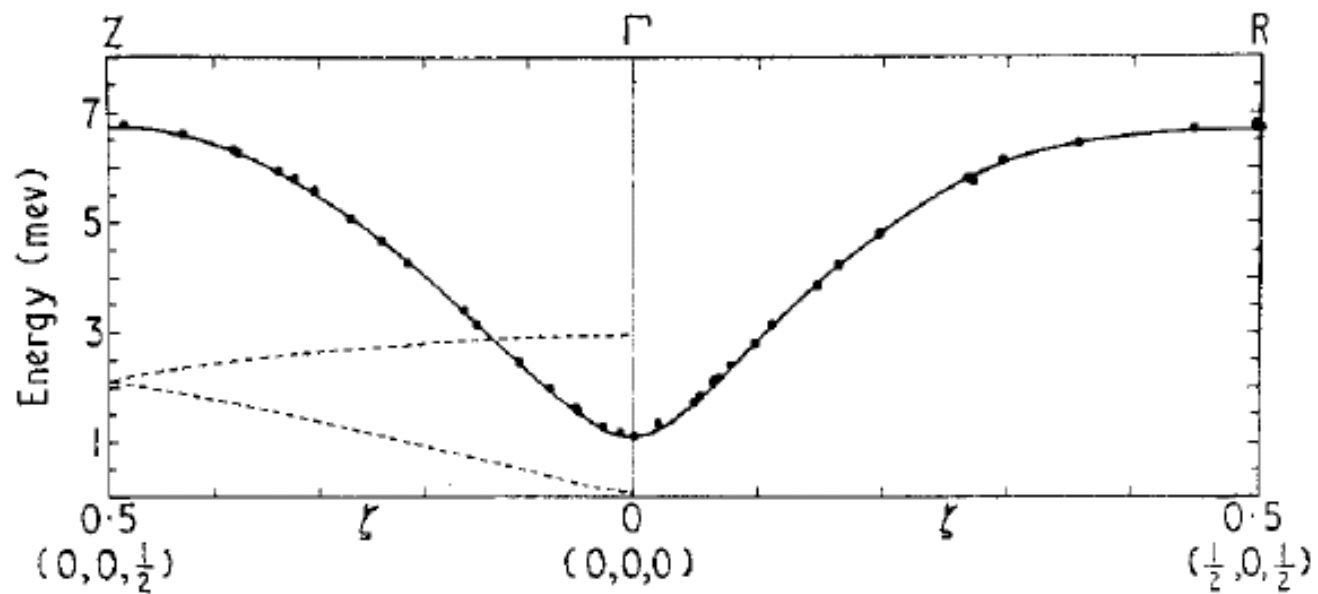
TbMnO₃ : l'axe b



On note une diminution du poids spectral des phonons à basse température, mais pas d'accident à T_N.

TbMnO₃ : l'axe c





1 meV = 8 cm⁻¹

NIKOTIN et al. *J. PHYS. C. (SOLID ST. PHYS.)*, 1969